## Lecture 3

Vector analysis and calculus
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From David K.Cheng CH2:
Vector Analysis

- Gradient of a Scalar Field
- Divergence of a Vector Field
- Divergence Theorem
- Curl of a Vector Field
- Stocke's Theorem

Differential operator $\nabla$ (Del operator):
$\nabla=\frac{\partial}{\partial x} \hat{a}_{x}+\frac{\partial}{\partial y} \hat{a}_{y}+\frac{\partial}{\partial z} \hat{a}_{z} \quad \nabla=\frac{\partial}{\partial \rho} \hat{a}_{\rho}+\frac{1}{\rho} \frac{\partial}{\partial \phi} \hat{a}_{\phi}+\frac{\partial}{\partial z} \hat{a}_{z} \quad \nabla=\frac{\partial}{\partial r} \hat{a}_{r}+\frac{1}{r} \frac{\partial}{\partial \theta} \hat{a}_{\theta}+\frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} \hat{a}_{\phi}$

It is a vector that its magnitude describes maximum rate of change of a scalar field and its direction describes the direction of this maximum rate of change at a given point in a given time.

$$
\nabla V=\frac{d V}{d n} \hat{a}_{n}=\frac{\partial V}{\partial x} \hat{a}_{x}+\frac{\partial V}{\partial y} \hat{a}_{y}+\frac{\partial V}{\partial z} \hat{a}_{z}
$$

Properties:

$$
\begin{aligned}
& \nabla(V+U)=\nabla V+\nabla U \\
& \nabla(V U)=(\nabla V) * U+(\nabla U) * V \\
& \nabla V^{n}=n V^{n-1} \nabla V
\end{aligned}
$$

EXAMPLE 2-16 The electrostatic field intensity $\mathbf{E}$ is derivable as the negative gradient of a scalar electric potential $V$; that is, $\mathbf{E}=-\nabla V$. Determine $\mathbf{E}$ at the point $(1,1,0)$ if
a) $V=V_{0} e^{-x} \sin \frac{\pi y}{4}$,

## Solution

a) $\mathbf{E}=-\left[\mathbf{a}_{x} \frac{\partial}{\partial x}+\mathbf{a}_{y} \frac{\partial}{\partial y}+\mathbf{a}_{z} \frac{\partial}{\partial z}\right] E_{0} e^{-x} \sin \frac{\pi y}{4}$

$$
=\left(\mathbf{a}_{x} \sin \frac{\pi y}{4}-\mathbf{a}_{y} \frac{\pi}{4} \cos \frac{\pi y}{4}\right) E_{0} e^{-x}
$$

Thus, $\mathbf{E}(1,1,0)=\left(\mathbf{a}_{x}-\mathbf{a}_{y} \frac{\pi}{4}\right) \frac{E_{0}}{\sqrt{2}}=\mathbf{a}_{E} E$,
where

$$
\begin{aligned}
E & =E_{0} \sqrt{\frac{1}{2}\left(1+\frac{\pi^{2}}{16}\right)} \\
\mathbf{a}_{E} & =\frac{1}{\sqrt{1+\left(\pi^{2} / 16\right)}}\left(\mathbf{a}_{x}-\mathbf{a}_{y} \frac{\pi}{4}\right) .
\end{aligned}
$$

$$
\begin{aligned}
& f=-x-y \\
& \nabla f=-a x-a y
\end{aligned}
$$


[x,y]=meshgrid(-5:1:5); [Vx,Vy]=gradient (-x-y) ; quiver(x,y,Vx,Vy) hold on; $[x, y]=m e s h g r i d(l i n s p a c e(-5,5,50))$;
$z=-x-y$; contour ( $x, y, z$ );hold off

$$
f(x, y)=9-x^{\wedge} 2-y^{\wedge} 2
$$

$$
\begin{aligned}
& \nabla f=-2 x \hat{a}_{x}-2 y \hat{a}_{y} \\
& \text { at }(2,-5) \rightarrow \nabla f=-4 \hat{a}_{x}+10 \hat{a}_{y}=10.7 \angle 111.8^{\circ} \\
& (-3,-1) \rightarrow \nabla f=6 \hat{a}_{x}+2 \hat{a}_{y}=7.2 \angle 18.4^{\circ} \\
& (-5,0) \rightarrow \nabla f=10 \hat{a}_{x}=10 \angle 0^{\circ} \\
& (2,0) \rightarrow \nabla f=-4 \hat{a}_{x}=4 \angle 180^{\circ}
\end{aligned}
$$



$[x, y]=m e s h g r i d(-5: 1: 5) ;[V x, V y]=$ gradient $\left(9-x . \wedge 2-y .{ }^{\wedge} 2\right) ; q u i v e r(x, y, V x, V y)$ hold on; $[x, y]=$ meshgrid(linspace $(-5,5,50)$ );
$z=9-x . \wedge 2-y .{ }^{\wedge} 2$;contour $(x, y, z)$;hold off

Divergence of a vector field $A$ is the net outward flux of $A$ per unit volume as the volume about the point tends to zero

$$
\operatorname{div} \vec{A} \approx \lim _{\Delta v \rightarrow 0} \frac{\oint_{s} \vec{A} \cdot d \vec{s}}{\Delta v}
$$

A divergence applied to a vector and result a scalar value that indicate the amount of vector field $A(r)$ that is Converging to, or diverging from ,a given point.


Positive
Divergence
source of flux at point $p$


Negative
Divergence
sink of flux at point $p$


Positive
Divergence
source of flux at point $p$


Zero
Divergence
Uniform field (solenoid field) or divergenceless field


Zero
Divergence
Nor source nor sink of flux at point $p$

Divergence of a Vector Field
If $d i v \bar{D}=\lim _{\Delta v \rightarrow 0} \frac{\oint \bar{D} . d \bar{s}}{\Delta v}$
prove
$\operatorname{div} \bar{D}=\left(\frac{\partial}{\partial x} \hat{a}_{x}+\frac{\partial}{\partial y} \hat{a}_{y}+\frac{\partial}{\partial z} \hat{a}_{z}\right) \cdot\left(D_{x} \hat{a}_{x}+D_{y} \hat{a}_{y}+D_{z} \hat{a}_{z}\right)=\frac{\partial D_{x}}{\partial x}+\frac{\partial D_{y}}{\partial y}+\frac{\partial D_{z}}{\partial z}$
We could use taylor expansion to get D at six faces of the cube
This term vanishes as $\Delta v$ reach 0 i.e x very close to a

$$
f(x)=f(a)+f^{\prime}(a) *(x-a)+f^{\prime \prime}(a) *(x-a)^{2}+\ldots \ldots .
$$

Taylor expansion: get value of function in required position by expansion of function in known position + derivative (changes Of function value from known position to desired position multiplied by distance between two positions

$$
\begin{aligned}
\oint_{S} \mathbf{D} & \cdot d \mathbf{S}=\int_{\text {front }}+\int_{\text {back }}+\int_{\text {left }}+\int_{\text {right }}+\int_{\text {top }}+\int_{\text {botton }} \\
\int_{\text {front }} & \doteq \mathbf{D}_{\text {front }} \cdot \Delta \mathbf{S}_{\text {front }} \\
& \doteq \mathbf{D}_{\text {front }} \cdot \Delta y \Delta z \mathbf{a}_{x} \\
& \doteq D_{\mathrm{x}, \text { front }} \Delta y \Delta z \\
D_{x, \text { front }} & \doteq D_{x 0}+\left.\frac{\Delta x}{2} \frac{\partial D_{x}}{\partial x}\right|_{x_{o}} \quad \text { (using taylor expansion) } \\
\int_{\text {front }} & \doteq\left(D_{x 0}+\frac{\Delta x}{2} \frac{\partial D_{x}}{\partial x}\right) \Delta y \Delta z
\end{aligned}
$$

$$
\begin{aligned}
\int_{\text {back }} & \doteq \mathbf{D}_{\text {back }} \cdot \Delta \mathbf{S}_{\text {back }} \\
& \doteq \mathbf{D}_{\text {back }} \cdot\left(-\Delta y \Delta z \mathbf{a}_{x}\right) \\
& \doteq-D_{x, \text { back }} \Delta y \Delta z
\end{aligned}
$$

$$
D_{x, \text { back }} \doteq D_{x 0}-\frac{\Delta x}{2} \frac{\partial D_{x}}{\partial x}
$$

$$
\int_{\text {back }} \doteq\left(-D_{x 0}+\frac{\Delta x}{2} \frac{\partial D_{x}}{\partial x}\right) \Delta y \Delta z
$$

$$
\int_{\text {front }}+\int_{\text {back }} \doteq \frac{\partial D_{x}}{\partial x} \Delta x \Delta y \Delta z
$$

similarly

$$
\begin{aligned}
& \int_{\text {right }}+\int_{\text {left }} \doteq \frac{\partial D_{y}}{\partial y} \Delta x \Delta y \Delta z \\
& \int_{\text {top }}+\int_{\text {bottom }} \doteq \frac{\partial D_{z}}{\partial z} \Delta x \Delta y \Delta z
\end{aligned}
$$

$$
\oint_{S} \mathbf{D} \cdot d \mathbf{S} \doteq\left(\frac{\partial D_{x}}{\partial x}+\frac{\partial D_{y}}{\partial y}+\frac{\partial D_{z}}{\partial z}\right) \Delta x \Delta y \Delta z
$$

or

$$
\oint_{S} \mathbf{D} \cdot d \mathbf{S}=Q \doteq\left(\frac{\partial D_{x}}{\partial x}+\frac{\partial D_{y}}{\partial y}+\frac{\partial D_{z}}{\partial z}\right) \Delta v
$$

$$
\lim _{\Delta v \rightarrow 0} \frac{\oint_{S} \mathbf{D} \cdot d \mathbf{S}}{\Delta v}=\lim _{\Delta v \rightarrow 0} \frac{Q}{\Delta v}=\left(\frac{\partial D_{x}}{\partial x}+\frac{\partial D_{y}}{\partial y}+\frac{\partial D_{z}}{\partial z}\right), \lim _{\Delta v \rightarrow 0} \frac{Q}{\Delta v}=\rho_{v}
$$

$$
\operatorname{div} \mathbf{D}=\left(\frac{\partial D_{x}}{\partial x}+\frac{\partial D_{y}}{\partial y}+\frac{\partial D_{z}}{\partial z}\right)
$$

## Example:

Given

$$
\bar{A}=x^{2} \quad \hat{a}_{x}+x y \hat{a}_{y}
$$

Find $\operatorname{div} \mathrm{A}$, plot A in arrow $\operatorname{div} \mathrm{A}$ in contour

$$
\nabla \cdot \bar{A}=2 x+x=3 x
$$

$[x, y]=$ meshgrid(-8:2:8,-8:2:8);
$\mathrm{Fx}=\mathrm{x} .{ }^{\wedge} 2$;
Fy = x.*y;
quiver (x,y,Fx,Fy)
D = divergence (x,y,Fx,Fy);
hold on
contour (x,y, D,'ShowText','on') hold off

EXAMPLE 2-17 Find the divergence of the position vector to an arbitrary point.
Solution We will find the solution in Cartesian as well as in spherical coordinates.
a) Cartesian coordinates. The expression for the position vector to an arbitrary point $(x, y, z)$ is

$$
\begin{gathered}
\overrightarrow{O P}=\mathbf{a}_{x} x+\mathbf{a}_{y} y+\mathbf{a}_{z} z . \\
\nabla \cdot(\overrightarrow{O P})=\frac{\partial x}{\partial x}+\frac{\partial y}{\partial y}+\frac{\partial z}{\partial z}=3 .
\end{gathered}
$$

b) Spherical coordinates. Here the position vector is simply

$$
\overrightarrow{O P}=\mathbf{a}_{R} R
$$

Its divergence in spherical coordinates $(R, \theta, \phi)$ can be obtained from

$$
\nabla \cdot \mathbf{A}=\frac{1}{R^{2}} \frac{\partial}{\partial R}\left(R^{2} A_{R}\right)+\frac{1}{R \sin \theta} \frac{\partial}{\partial \theta}\left(A_{\theta} \sin \theta\right)+\frac{1}{R \sin \theta} \frac{\partial A_{\phi}}{\partial \phi} .
$$

$$
\nabla \cdot(\overrightarrow{O P})=3
$$

## Divergence theorem

We define divergence as the net outward flux per unit volume
So the volume integral of the divergence of a vector field equals the total outward flux of the vector through the surface that bounds the volume

$$
\int_{v} \nabla \cdot \vec{A} d v=\oint_{s} \bar{A} \cdot d \bar{s}
$$

EXAMPLE 2-19 Given $\mathbf{A}=\mathbf{a}_{x} x^{2}+\mathbf{a}_{y} x y+\mathbf{a}_{z} y z$, verify the divergence theorem over a cube one unit on each side. The cube is situated in the first octant of the Cartesian coordinate system with one corner at the origin.

Solution Refer to Fig. 2-28. We first evaluate the surface integral over the six faces.

1. Front face: $x=1, d \mathbf{s}=\mathbf{a}_{x} d y d z$;

$$
\int_{\substack{\text { front } \\ \text { face }}} \mathbf{A} \cdot d \mathbf{s}=\int_{0}^{1} \int_{0}^{1} d y d z=1
$$

2. Back face: $x=0, d \mathbf{s}=-\mathbf{a}_{x} d y d z$;

$$
\int_{\text {fack }} \mathbf{A} \cdot d \mathbf{s}=0
$$

3. Left face: $y=0, d \mathbf{s}=-\mathbf{a}_{y} d x d z$;

$$
\int_{\substack{\text { left } \\ \text { face }}} \mathbf{A} \cdot d \mathbf{s}=0
$$


$\mathbf{A}=\mathbf{a}_{x} x^{2}+\mathbf{a}_{v} x y+\mathbf{a}_{z} y z$,
4. Right face: $y=1, d \mathbf{s}=\mathbf{a}_{y} d x d z$;

$$
\int_{\substack{\text { right } \\ \text { face }}} \mathbf{A} \cdot d \mathbf{s}=\int_{0}^{1} \int_{0}^{1} x d x d z=\frac{1}{2} .
$$

5. Top face: $z=1, d \mathbf{s}=\mathbf{a}_{z} d x d y$;

$$
\int_{\substack{\text { top } \\ \text { face }}} \mathbf{A} \cdot d \mathbf{s}=\int_{0}^{1} \int_{0}^{1} y d x d y=\frac{1}{2} .
$$

6. Bottom face: $z=0, d \mathbf{s}=-\mathbf{a}_{z} d x d y$;

$$
\int_{\substack{\text { bottom } \\ \text { face }}} \mathbf{A} \cdot d \mathbf{s}=0 .
$$

Adding the above six values, we have

$$
\oint_{S} \mathbf{A} \cdot d \mathbf{s}=1+0+0+\frac{1}{2}+\frac{1}{2}+0=2
$$

Now the divergence of $\mathbf{A}$ is

$$
\boldsymbol{\nabla} \cdot \mathbf{A}=\frac{\partial}{\partial x}\left(x^{2}\right)+\frac{\partial}{\partial y}(x y)+\frac{\partial}{\partial z}(y z)=3 x+y
$$

Hence,

$$
\int_{V} \boldsymbol{\nabla} \cdot \mathbf{A} d v=\int_{0}^{1} \int_{0}^{1} \int_{0}^{1}(3 x+y) d x d y d z=2
$$

which is the same as the result of the closed surface integral in (2-120). The divergence theorem is therefore verified.

## Curl of a Vector Field

Curl of a vector field $\underline{H}(\nabla \times \underline{H})$ is a vector $(\underline{J})$ whose magnitude is the maximum net circulating of H per unit area as area tends to zero and whose direction is the normal direction of the area (when the area is oriented to make the net circulation Maximum)

$$
\text { Used in ampere law: } \nabla \times \underline{H}=\underline{\jmath} \text { or } \quad \oint \bar{H} \cdot d \bar{l}=I_{\text {enc }}=\bar{J} \cdot d \bar{s}
$$

$$
\nabla \times \bar{H}=\lim _{\Delta S_{n} \rightarrow 0} \frac{\oint_{c} \bar{H} \cdot d \bar{l}}{\Delta S_{n}}
$$

The curl of a vector field $A$ at a point $P$ may indicates how much the field rotate around $P$ or how much field vary In direction normal to its flow direction (if we put pin it cause it to rotate)

$\nabla X A>0 \quad \nabla X A<0$
$\nabla X A=0$
$\oint_{c} \bar{H} \cdot d \bar{l}=(H . \Delta L)_{1-2}+(H . \Delta L)_{2-3}+(H . \Delta L)_{3-4}+(H . \Delta L)_{4-1}$
$(\bar{H} . \Delta \bar{L})_{1-2}=H_{y, 1-2} \Delta y=\left(H_{y 0}+\frac{\Delta x}{2} \frac{\partial H_{y}}{\partial x}\right) \Delta y$
$(\bar{H} . \Delta \bar{L})_{3-4}=H_{y, 3-4}(-\Delta y)=\left(H_{y 0}-\frac{\Delta x}{2} \frac{\partial H_{y}}{\partial x}\right)(-\Delta y)$
$(\bar{H} \cdot \Delta \bar{L})_{1-2}+(\bar{H} \cdot \Delta \bar{L})_{3-4}=\Delta x \Delta y \frac{\partial H_{y}}{\partial x}$
similarly
$(\bar{H} . \Delta \bar{L})_{2-3}=H_{x, 2-3}(-\Delta x)=-\left(H_{x 0}+\frac{\Delta y}{2} \frac{\partial H_{x}}{\partial y}\right) \Delta x$
$\mathbf{H}=\mathbf{H}_{0}=H_{x 0} \mathbf{a}_{x}+H_{y 0} \mathbf{a}_{y}+H_{z 0} \mathbf{a}_{z}$
$(\bar{H} . \Delta \bar{L})_{4-1}=H_{x, 4-1}(\Delta x)=\left(H_{x 0}-\frac{\Delta y}{2} \frac{\partial H_{x}}{\partial y}\right)(\Delta x)$
$(\bar{H} \cdot \Delta \bar{L})_{2-3}+(\bar{H} \cdot \Delta \bar{L})_{4-1}=-\Delta x \Delta y \frac{\partial H_{x}}{\partial y}$
$\oint_{c} \bar{H} \cdot d \bar{l}=\Delta x \Delta y\left(\frac{\partial H_{y}}{\partial x}-\frac{\partial H_{x}}{\partial y}\right)=J_{z} \Delta x \Delta y \rightarrow \lim _{\Delta x \Delta y \rightarrow 0} \frac{\oint_{c} \bar{H} \cdot d \bar{l}}{\Delta x \Delta y}=J_{z}$

Choosing closed paths at yz plane and xz plane and do analogous processes lead to expressions for the $y$ and $z$ components Of the current density

$$
\lim _{\Delta y, \Delta z \rightarrow 0} \frac{\oint \mathbf{H} \cdot d \mathbf{L}}{\Delta y \Delta z}=\frac{\partial H_{z}}{\partial y}-\frac{\partial H_{y}}{\partial z}=J_{x}
$$

$$
\lim _{\Delta z, \Delta x \rightarrow 0} \frac{\oint \mathbf{H} \cdot d \mathbf{L}}{\Delta z \Delta x}=\frac{\partial H_{x}}{\partial z}-\frac{\partial H_{z}}{\partial x}=J_{y}
$$

$$
\begin{aligned}
& \lim _{\Delta S_{N} \rightarrow 0} \frac{\oint \mathbf{H} \cdot d \mathbf{L}}{\Delta S_{N}}=\left(\frac{\partial H_{z}}{\partial y}-\frac{\partial H_{y}}{\partial z}\right) \mathbf{a}_{x}+\left(\frac{\partial H_{x}}{\partial z}-\frac{\partial H_{z}}{\partial x}\right) \mathbf{a}_{y}+\left(\frac{\partial H_{y}}{\partial x}-\frac{\partial H_{x}}{\partial y}\right) \mathbf{a}_{z} \\
& \quad \operatorname{curl} \mathbf{H}=\nabla \times \mathbf{H}=\left(\frac{\partial H_{z}}{\partial y}-\frac{\partial H_{y}}{\partial z}\right) \mathbf{a}_{x}+\left(\frac{\partial H_{x}}{\partial z}-\frac{\partial H_{z}}{\partial x}\right) \mathbf{a}_{y}+\left(\frac{\partial H_{y}}{\partial x}-\frac{\partial H_{x}}{\partial y}\right) \mathbf{a}_{z}=\mathbf{J}
\end{aligned}
$$

$$
\nabla \times \mathbf{H}=\mathbf{J} \quad \text { Ampère's circuital law }
$$

## Example 8.2

As an example of the evaluation of curl $\mathbf{H}$ from the definition and of the evaluation of another line integral, let us suppose that $\mathbf{H}=0.2 z^{2} \mathbf{a}_{x}$ for $z>0$, and $\mathbf{H}=0$ elsewhere, as shown in Fig. 8.15. Calculate $\oint \mathbf{H} \cdot d \mathbf{L}$ about a square path with side $d$, centered at $\left(0,0, z_{1}\right)$ in the $y=0$ plane where $z_{1}>2 d$.

Solution. We evaluate the line integral of $\mathbf{H}$ along the four segments, beginning at the top:

$$
\begin{array}{lrl}
\qquad \begin{aligned}
\oint \mathbf{H} \cdot d \mathbf{L} & =0.2\left(z_{1}+\frac{1}{2} d\right)^{2} d+0-0.2\left(z_{1}-\frac{1}{2} d\right)^{2} d+0 \\
& =0.4 z_{1} d^{2}
\end{aligned} & \text { or }=d x d z \frac{\partial \boldsymbol{H} \boldsymbol{x}}{\partial z}
\end{array} \text { In the limit as the area approaches zero, we find } \quad ~ a s ~ \nabla \times \bar{H}=\lim _{\Delta s_{n} \rightarrow 0} \frac{\oint \bar{H} \cdot d \bar{l}}{\Delta S_{n}} .
$$



$$
(\nabla \times \mathbf{H})_{y}=\lim _{d \rightarrow 0} \frac{\oint \mathbf{H} \cdot d \mathbf{L}}{d^{2}}=\lim _{d \rightarrow 0} \frac{0.4 z_{1} d^{2}}{d^{2}}=0.4 z_{1}
$$

$$
\nabla \times \mathbf{H}=\left|\begin{array}{ccc}
\mathbf{a}_{x} & \mathbf{a}_{y} & \mathbf{a}_{z} \\
\frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\
0.2 z^{2} & 0 & 0
\end{array}\right|=\frac{\partial}{\partial z}\left(0.2 z^{2}\right) \mathbf{a}_{y}=0.4 z \mathbf{a}_{y}
$$

which checks with the result above when $z=z_{1}$.


## Another example

$$
\begin{aligned}
\vec{f} & =\sin (y) \hat{a}_{x}+\sin (x) \hat{a}_{y} \\
\nabla & \times \vec{f}=\frac{\partial}{\partial x}(\sin (x))-\frac{\partial}{\partial y}(\sin (y)) \quad \hat{a}_{z} \\
& =\cos (x)-\cos (y) \quad \hat{a}_{z}
\end{aligned}
$$



$$
\lim _{\Delta s_{j} \rightarrow 0} \sum_{j=1}^{N}(\nabla \times \bar{H})_{j} \Delta \bar{s}_{j}=\int_{s} \nabla \times \bar{H} \cdot d \bar{s}
$$

## Stocks' Theorem:

The surface integral of the curl of a vector field over an open surface is equal to the closed line integral of the vector along the contour bounding the surface
$\oint_{s}^{\text {Note: }} \nabla \times \bar{H} \cdot d \bar{s}=0$


## FIGURE 8.16

The sum of the closed line integrals about the perimeter of every $\Delta S$ is the same as the closed line integral about the perimeter of $S$ because of cancellation on every interior path.

For any closed surface no surface bounding contour c exist so curl over any closed surface=0

## Example

A numerical example may help to illustrate the geometry involved in Stokes' theorem. Consider the portion of a sphere shown in Fig. 8.17. The surface is specified by $r=4$, $0 \leq \theta \leq 0.1 \pi, 0 \leq \phi \leq 0.3 \pi$, and the closed path forming its perimeter is composed of three circular arcs. We are given the field $\mathbf{H}=6 r \sin \phi \mathbf{a}_{r}+18 r \sin \theta \cos \phi \mathbf{a}_{\phi}$ and are asked to evaluate each side of Stokes' theorem.
Solution. The first path segment is described in spherical coordinates by $r=4$, $0 \leq \theta \leq 0.1 \pi, \phi=0$; the second one by $r=4, \theta=0.1 \pi, 0 \leq \phi \leq 0.3 \pi$; and the third by $r=4,0 \leq \theta \leq 0.1 \pi, \phi=0.3 \pi . \quad d \mathbf{L}=d r \mathbf{a}_{r}+r d \theta \mathbf{a}_{\theta}+r \sin \theta d \phi \mathbf{a}_{\phi}$

$$
\begin{gathered}
H_{\theta}=0, \\
\oint \mathbf{H} \cdot d \mathbf{L}=\int_{1} H_{\theta} r d \theta+\int_{2} H_{\phi} r \sin \theta d \phi+\int_{3} H_{\phi} r d \theta \\
\oint \mathbf{H} \cdot d \mathbf{L}=\int_{0}^{0.3 \pi}[18(4) \sin 0.1 \pi \cos \phi] 4 \sin 0.1 \pi d \phi=288 \sin ^{2} 0.1 \pi \sin 0.3 \pi=22.2 \mathrm{~A}
\end{gathered}
$$

$$
\nabla \times \mathbf{H}=\frac{1}{r \sin \theta}(36 r \sin \theta \cos \theta \cos \phi) \mathbf{a}_{r}+\frac{1}{r}\left(\frac{1}{\sin \theta} 6 r \cos \phi-36 r \sin \theta \cos \phi\right) \mathbf{a}_{\theta}
$$

Since $d \mathbf{S}=r^{2} \sin \theta d \theta d \phi \mathbf{a}_{r}$, the integral is

$$
\begin{aligned}
\int_{S}(\nabla \times \mathbf{H}) \cdot d \mathbf{S} & =\int_{0}^{0.3 \pi} \int_{0}^{0.1 \pi}(36 \cos \theta \cos \phi) 16 \sin \theta d \theta d \phi \\
& =\left.\int_{0}^{0.3 \pi} 576\left(\frac{1}{2} \sin ^{2} \theta\right)\right|_{0} ^{0.1 \pi} \cos \phi d \phi \\
& =288 \sin ^{2} 0.1 \pi \sin 0.3 \pi=22.2 \mathrm{~A}
\end{aligned}
$$

Assignment (Due date 8/11/2020)
-research on: electromagnetic train

Project:
-electromagnetic breaking system
-RFID Access control Using Arduino Required:
-literature review
-theory of operation
-design circuit
Site containing different project ideas in various fields:
https://courses.engr.illinois.edu/ece445/projects.asp

